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Noncovalent dyads of lanthanide nitride cluster fullerenes Ln₃N@C₈₀ and bisphthalocyanines LnPc₂: A DFT Study

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Phthalocyanines (Pcs) Azomethine nitrogen Peripheral (or b)-substituents (or a)-substituents MPc (M = 3d metal transition) Background D WPcs MPc (M = rare earth elements, including Y and lanthanides); Ln = lanthanides.

Application: Spintronic devices

Single molecule magnets (SMMs)

Saccelrys°

MATERIALS STUDIO

Module **DMol**³

Carbon nanomaterials + MPc₂

Noncovalent interactions between LnPc₂ and carbon nanotubes as well as graphene (with defects), have been investigated using density functional theory (DFT). The most relevant conclusion from the published studies is that the interaction of GdPc₂ with zigzag CNTs ^[2] and graphene with 5665 defects ^[3] leads to changes in spin orientation compared to the isolated components.

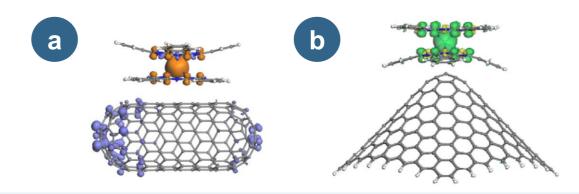


Fig.2. Spin density plots for (**a**) $GdPc_2+SWCNTs$ and (**b**) $GdPc_2+G-5665$ (Isosurfaces at 0.01 a.u.). [2,3]

a) Violet and orange lobes represent spin-up and spin-down electrons.b) Green and yellow lobes represent spin-up and spin-down electrons

Nitride cluster fullerene (Ln₃N@C₈₀; NCFs)

What changes occurs in HOMO, LUMO, and spin density distribution of $LnPc_2$ (Ln = La, Ce, Gd, Lu) upon interaction with $Ln_3N@C_{80}$, compared to the isolated $LnPc_2$ systems?



Aim

 To analyze the geometric and electronic property changes of LnPc₂ upon noncovalent interaction with Ln₃N@C₈₀ using DFT.



Results and discussion

Computational method

DFT

Functional: general gradient approximation (**GGA**) functional by Perdew-Burke-Ernzerhof (**PBE**) in combination with a long-range dispersion correction by Grimme (**PBE-D**).

Fig.1. Structure of (a) monophthalocyanines and (b) bisphthalocyanine. [1]

Different basis sets (DN, DND, DNP) were tested; structures were selected based on agreement with XRD data. DN was used to simulate the dyads.

Optimization of isolated bisphthalocyanines and nitride cluster fullerenes

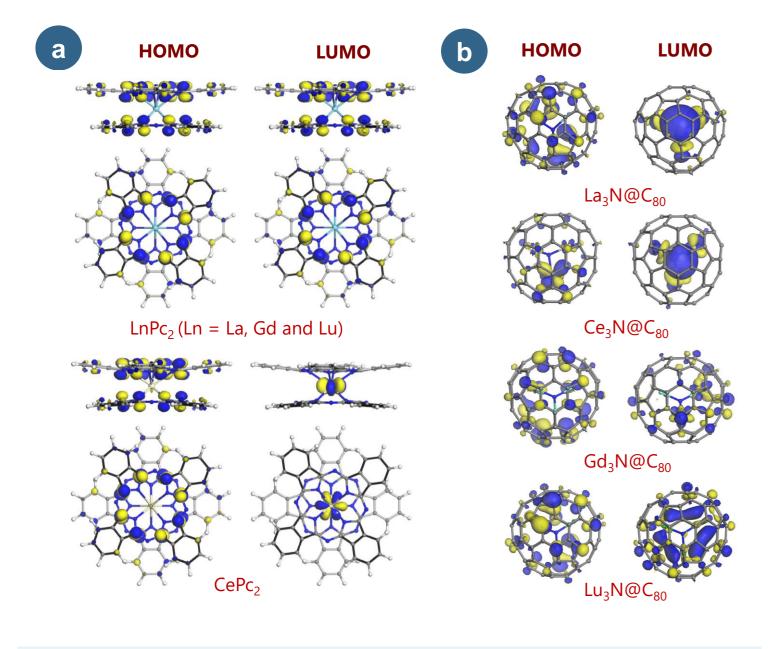


Fig. 3. Frontier orbital distribution (HOMO and LUMO; isosurfaces at 0.03 a.u.) for isolated (a) bisphthalocyanine and (b) $Ln_3N@C_{80}$.

Table 1. HOMO-LUMO gap energies (in eV) for isolated components Ln₃N@C₈₀ and LnPc₂.

	System	E _{gap} (eV)
	La ₃ N@C ₈₀	1.371
TIME TO LITER	$Ce_3N@C_{80}$	0.862
	$Gd_3N@C_{80}$	1.457
	$Lu_3N@C_{80}$	1.528
URSONA NICONAL AUTORONA O METURO	LaPc ₂	0.131
POR AT POWOUL	CePc ₂	0.196
Cure State	$GdPc_2$	0.130
	LuPc ₂	0.137
JNAW		

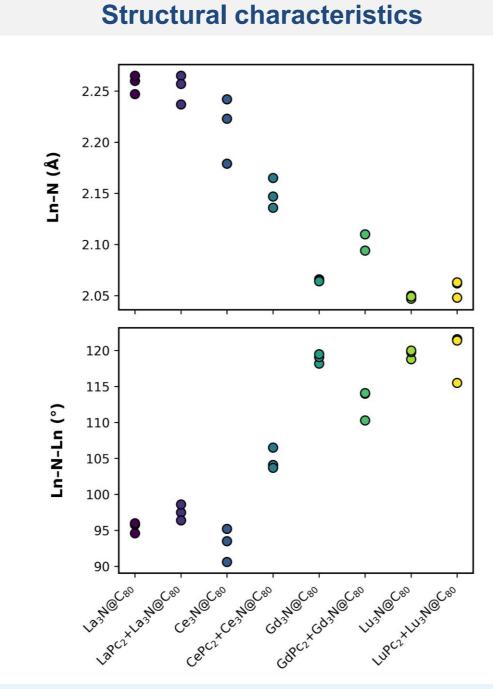


Fig. 4. Comparison of the Ln-N bond lengths (Å; top) and Ln-N-Ln angles (°; bottom) in the isolated lanthanide nitride cluster $Ln_3N@C_{80}$ (Ln = La, Ce, Gd, and Lu) and in the noncovalent $LnPc_2 + Ln_3N@C_{80}$ dyads.

The changes in spin density distribution are evident in the dyads containing Ce and Gd atoms, contrary to their La and Lu-derived counterparts. The interaction of $Ce_3N@C_{80}$ and $Gd_3N@C_{80}$ with $CePc_2$ and $GdPc_2$, respectively, causes redistribution of the spin density, with changes in the orientation of spin-up and spin-down electrons in the encapsulated Ce_3N and Gd_3N clusters.

Acknowledgements

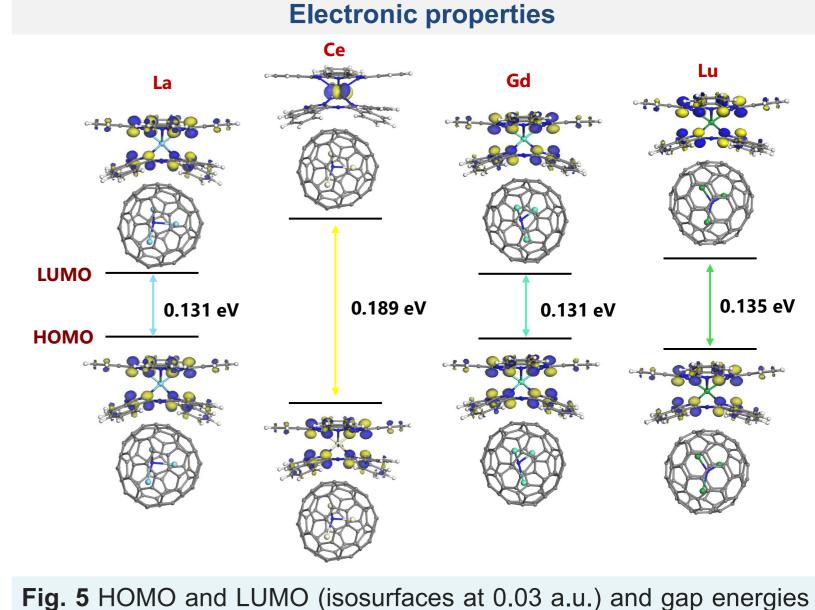
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References

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LnPc₂ + La₃N@C₈₀ dyads



for noncovalent $LnPc_2 + Ln_3N@C_{80}$ dyads.

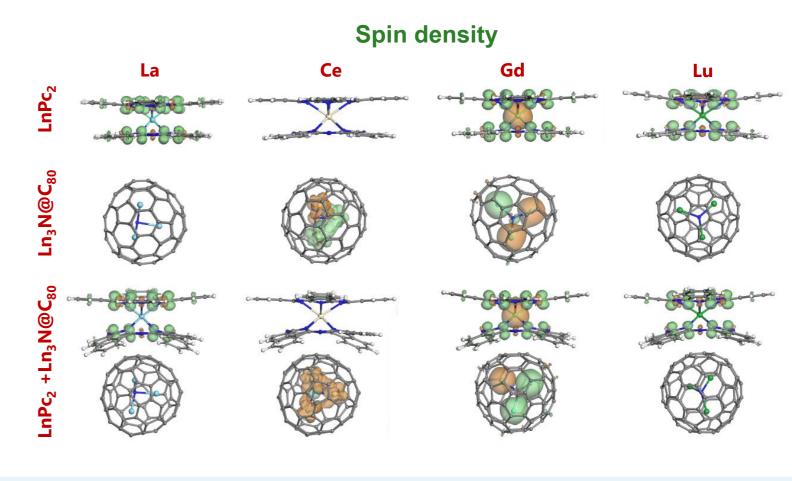


Fig. 6. Comparison of spin density patterns (isosurfaces at 0.01 a. u.) for isolated $LnPc_2$ and $Ln_3N@C_{80}$ molecules with those for their noncovalent dyads $LnPc_2 + Ln_3N@C_{80}$. The green and orange lobes correspond to spin-up and spin-down electrons, respectively.

Conclusions

- Using a smaller basis set (DN) proved effective for lanthanide species, with no convergence issue
- Noncovalent LnPc₂ + Ln₃N@C₈₀ interactions induce structural changes in the encapsulated Ln₃N cluster, favoring planar or pyramidal geometries.
- Energy gap values are closer to those of isolated bisphthalocyanines, indicating that the LnPc₂ component governs the reactivity of the dyad.