



Gapless Superfluidity and Neutron Star Cooling [†]

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Abstract: The presence of currents in the interior of cold neutron stars can lead to a state in which nucleons remain superfluid while the quasiparticle energy spectrum has no gap. We show within the self-consistent time-dependent nuclear energy density functional theory that the nucleon specific heat can then become comparable to that in the normal phase in contrast to the classical BCS result in the absence of superflows. This dynamical gapless superfluid state has important implications for the cooling of neutron stars.

Keywords: neutron star; superfluidity; cooling; pairing gap; specific heat

1. Introduction

Produced during gravitational-core collapse supernova explosions with initial temperatures as high as $\sim 10^{11} - 10^{12}$ K, neutron stars cool down to temperatures $\sim 10^9$ K within a few days [1]. The very dense matter in their interior is expected to undergo various quantum phase transitions analogous to those observed in terrestrial laboratories [2]. Similarly to electrons in conventional terrestrial superconductors, free neutrons in the inner crust and the outer core of neutron stars are predicted to form a Bardeen-Cooper-Schrieffer (BCS) condensate of 1S_0 Cooper pairs. Nuclear superfluidity has found support from the rapid decline of luminosity of the Cassiopeia A remnant [3–9] and has been corroborated by radio-timing observations of frequency glitches in numerous pulsars [10] interpreted as global readjustments of the rotational motions of the neutron and proton superfluids induced by the unpinning of quantized neutron superfluid vortices [11,12] (see, e.g., [13] for a review).

Despite the importance of the superfluid dynamics for interpreting these latter astrophysical phenomena, most microscopic calculations of the nuclear pairing properties have been carried out so far for static situations (see, e.g., [14] for a recent review). We have recently studied the dynamics of hot neutron-proton superfluid mixtures within the self-consistent time-dependent nuclear energy density functional theory [15]. In application to neutron stars, we have computed 1S_0 neutron and proton pairing gaps in the homogeneous core in the presence of arbitrary currents and we have determined the mutual neutron-proton entrainment coupling coefficients [16]. We have also shown within the same framework that there exists a dynamical “gapless” state in which nuclear superfluidity is not destroyed even though the energy spectrum of quasiparticle excitations exhibits no gap [17]. As will be shown in Section 2, the absence of an energy gap leads to a nucleon specific heat that is very different from that in the classical BCS state (in the absence of superflows). The implications for the cooling of neutron stars will be discussed in Section 3.

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2. Gapless Superfluidity

2.1. Order Parameter

We have previously studied the nuclear superfluidity at finite temperatures and in the presence of superflows under the self-consistent time dependent nuclear energy density functional framework [15,16]. The behavior of the order parameter Δ_q (with $q = n, p$ for neutron and proton respectively) has been found to be universal after introducing some effective superfluid velocity \mathbb{V}_q and proper rescaling (see Figure 1): $\Delta_q^{(0)}$ is the order parameter at zero temperature and in the absence of superflows, $T_{cq}^{(0)} = e\gamma\Delta_q^{(0)}/\pi \simeq 0.567\Delta_q^{(0)}$ is the critical temperature above which superfluidity is destroyed and $\mathbb{V}_{Lq} = \Delta_q^{(0)}/(\hbar k_{Fq})$ (with k_{Fq} , the Fermi wave vector) is Landau's velocity (derived via the eponymous criterion [18,19] adapted to the context of strongly interacting nuclear superfluid mixtures).

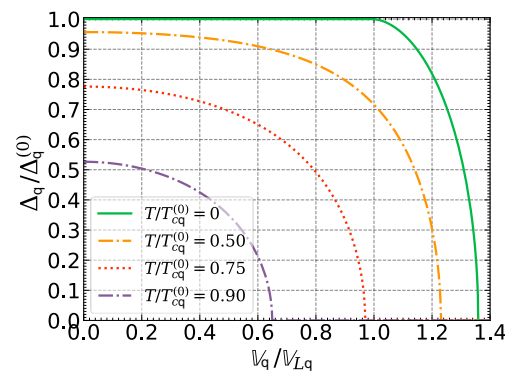


Figure 1. 1S_0 nucleon pairing gap Δ_q normalized to its value $\Delta_q^{(0)}$ at zero temperature in the absence of superflow as a function of the corresponding effective superfluid velocity \mathbb{V}_q expressed in terms of Landau's velocity \mathbb{V}_{Lq} for different temperatures ($T_{cq}^{(0)}$ being the critical temperature for the superfluid at rest). See text for details.

2.2. Gapless State and Specific Heat

Focusing on low temperatures, $T \ll T_{cq}^{(0)}$ (relevant for mature neutron stars), we have shown that the energy gap in the quasiparticle density of state \mathcal{D}_q shrinks with increasing effective superfluid velocity \mathbb{V}_q and disappears at Landau's velocity \mathbb{V}_{Lq} (see Figure 2). However, the order parameter Δ_q remains finite: the superfluid enters a gapless state, which persists until the critical velocity $\mathbb{V}_{cq}^{(0)} = e\mathbb{V}_{Lq}/2 \simeq 1.36\mathbb{V}_{Lq}$ is reached [17].

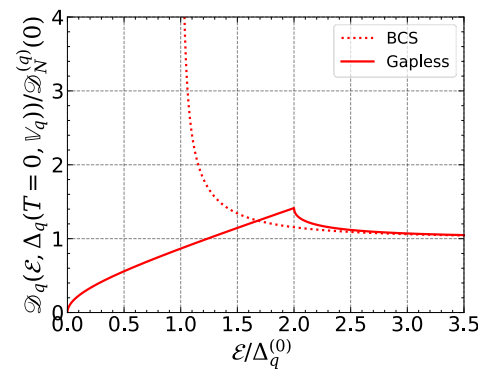


Figure 2. Quasiparticle density of states at low temperatures (normalized by the one in the normal phase $\mathcal{D}_N^{(q)}(0)$) in the BCS state ($\mathbb{V}_q = 0$) and in the gapless state ($\mathbb{V}_q = \mathbb{V}_{Lq}$).

One of the immediate consequences of gapless superfluidity is the modification of thermal properties such as the specific heat $c_V^{(q)}$. While the classical BCS state leads to an exponentially suppressed specific heat at low temperatures (see, e.g., [20]),

$$c_V^{(q)}(T \ll T_{cq}^{(0)}, \mathbb{V}_q = 0) \approx \frac{3\sqrt{2}}{\pi^{3/2}} \left(\frac{T_{cq}^{(0)} \pi}{T} e^\gamma \right) \exp\left(\frac{T_{cq}^{(0)} \pi}{T} e^\gamma \right) c_N^{(q)}(T), \tag{1}$$

$c_N^{(q)}$ being the corresponding specific heat in the normal phase and $\gamma \simeq 0.577216$ denoting the Euler-Mascheroni constant, the specific heat in gapless state is comparable to that in the normal phase, and is approximately given by [17]

$$c_V^{(q)}(T \ll T_{cq}^{(0)}, \mathbb{V}_q > \mathbb{V}_{Lq}) \approx \sqrt{1 - \left(\frac{\Delta_q(\mathbb{V}_q) \mathbb{V}_{Lq}}{\Delta_q^{(0)} \mathbb{V}_q} \right)^2} c_N^{(q)}(T), \tag{2}$$

and Δ_q is accurately given by the following interpolation [16]:

$$\frac{\Delta_q(\mathbb{V}_q > \mathbb{V}_{Lq})}{\Delta_q^{(0)}} = 0.5081 \sqrt{1 - \frac{2\mathbb{V}_q}{e\mathbb{V}_{Lq}} \left(2.437 \frac{\mathbb{V}_q}{\mathbb{V}_{Lq}} - 4.443 \sqrt{\frac{\mathbb{V}_{Lq}}{\mathbb{V}_q}} + 5.842 \right)}. \tag{3}$$

The ratio between the specific heat with its associated specific heat in the normal phase is therefore an increasing universal function of $\mathbb{V}_q/\mathbb{V}_{Lq}$.

3. Astrophysical Consequences

Gapless superfluidity has important implications for the cooling of neutron stars, and in particular for the interpretation of the thermal emission from quasipersistent soft X-ray transients [21]. These binary systems consist of a neutron star whose crust is sporadically heated due to mass-transfer from a low-mass stellar companion for a long period of time (from years to decades) before entering a cooling phase when the accretion stops. The thermal relaxation has been observed for several sources up to about 10^4 days after outbursts [22] and is governed by the diffusion of heat in the inner crust of neutron stars, which is made of ions, free electrons and superfluid neutrons (see, e.g., [23]).

Let us recall that the typical thermal timescale of a crustal layer, delimited by the radial coordinates r_{\min} and r_{\max} , is given by [24]

$$\tau \approx \frac{1}{4} \left(\int_{r_{\min}}^{r_{\max}} dr \sqrt{\frac{c_V}{\kappa}} \right)^2, \tag{4}$$

with κ the thermal conductivity and c_V the total specific heat. So far, the neutron contribution to the crustal specific heat has been generally thought to be negligible, assuming superfluid neutrons are in the classical BCS state. In this case, the thermal relaxation is expected to be much faster if neutrons are superfluid since τ is shorter according to equation (4). To explain the observed late time cooling of some sources, some authors proposed that neutrons are not superfluid in the deepest region of the crust [25,26]. However, this interpretation has been recently ruled out by microscopic calculations [27,28].

Astrophysical observations and nuclear physics can be reconciled by allowing neutrons to be in the gapless superfluid state [21]. Indeed, the neutron specific heat is then strongly enhanced compared to that in the BCS state and can now dominate the electronic and ionic contributions even when considering realistic nuclear pairing properties [27,28]. The end result is a delayed thermal relaxation of the neutron star crust, as observed.

4. Conclusions

Considering the dynamics of nuclear superfluidity in the framework of the self-consistent time-dependent nuclear energy-density functional theory [15,16], we have shown that the energy spectrum of quasiparticle excitations exhibits no gap while the order parameter Δ_q remains finite at low temperatures for effective superfluid velocities \mathbb{V}_q exceeding Landau's velocity \mathbb{V}_{Lq} but lower than the critical velocity $\mathbb{V}_{cq}^{(0)}$.

Contrary to the classical BCS state, the gapless state is characterized by a very large specific heat comparable to that in the normal phase. This specific heat at low temperatures is shown to be a universal function of V_q/V_{Lq} [17].

Focusing on the inner crust of neutron stars, we have shown that the drastic increase of the neutron specific heat can solve the apparent contradiction between the observed late time cooling of quasipersistent soft X-ray transients and microscopic nuclear pairing calculations [21].

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