



On cosmological inflation in Palatini $F(R, \varphi)$ gravity

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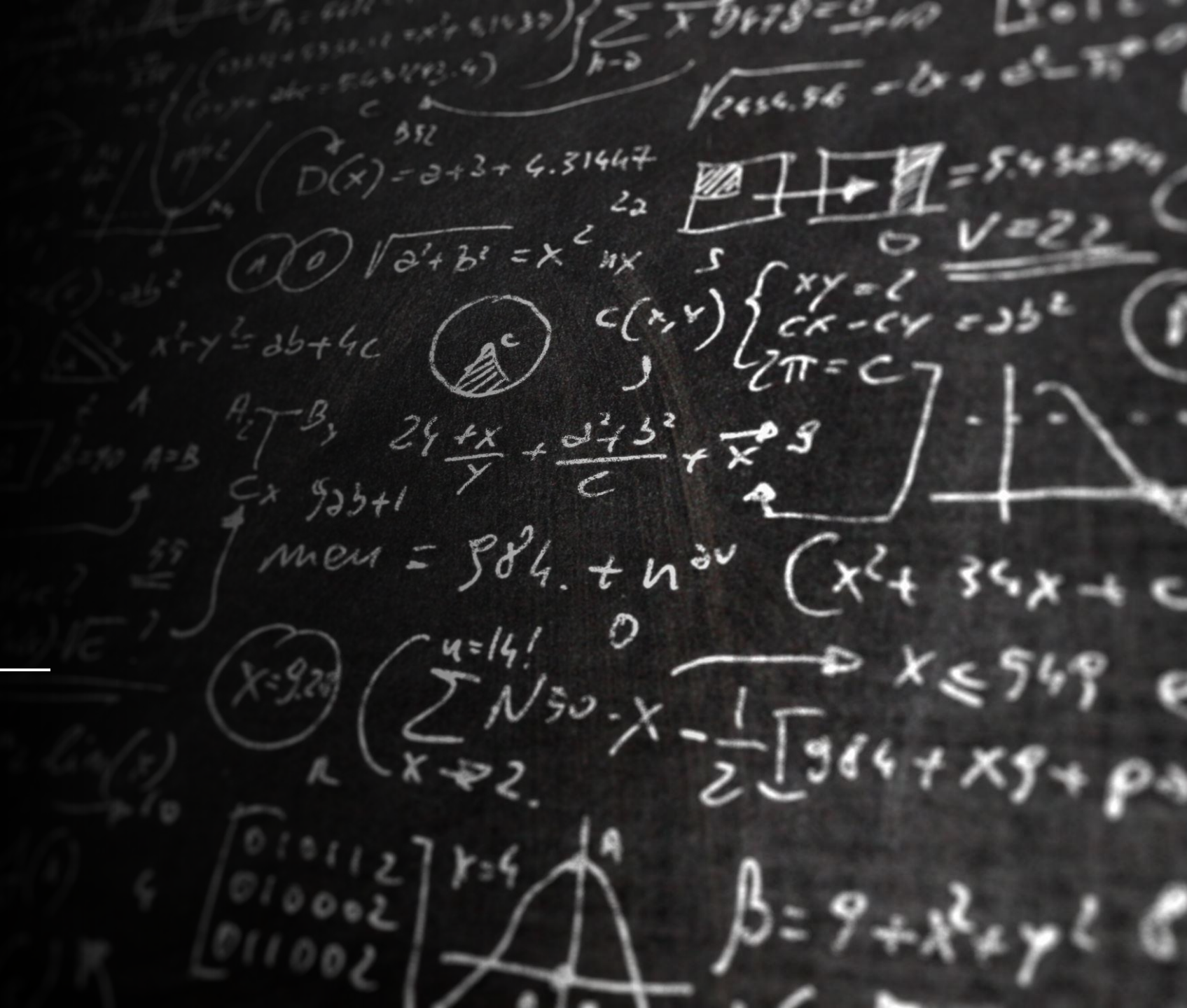




Contents:

- Why $F(R,\phi)$ gravity?
- Analytical treatment.
- Cosmological inflation in $F(R,\phi)$
- Case study: cosine and exponential potentials

Motivation for
 $F(R, \phi)$
theories of
gravity



Historical Motivation:

Scientific Curiosity



Eddington, A. S.,
1923, The
Mathematical
Theory of
Relativity
(Cambridge
University Press,
Cambridge).

Historical Motivation:

Scientific Curiosity



Eddington, A. S., 1923, The
Mathematical Theory of
Relativity (Cambridge
University Press, Cambridge).

Historical Motivation:

Scientific Curiosity



H.Weyl, A New Extension of
Relativity Theory, Annalen
Phys. 59(1919), 101-133
doi:10.1002/andp.191936410
02



Contemporary Motivation

Example: Explaining the Current accelerated expansion

$$S = S_{g,\varphi} + S_{\varphi}$$

The Model

$$S = \int d^4x \sqrt{-g} [F(R, \varphi)] + \int d^4x \sqrt{-g} \left[\frac{1}{2} g^{\mu\nu} \partial_{\mu}\varphi \partial_{\nu}\varphi - V(\varphi) \right]$$

By introducing an auxiliary scalar field, one can re-express the last theory as Brans-Dicke theory

$$S = \int d^4x \sqrt{-g} [(\chi + f(\varphi))R - \frac{1}{2} g^{\mu\nu} \partial_{\mu}\varphi \partial_{\nu}\varphi - V(\varphi) - V_{eff}(\chi)]$$

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By Applying a conformal transformation as:

$$\tilde{g}_{\mu\nu} = \Omega^2 g_{\mu\nu} \qquad \Omega^2 = (\chi + f(\varphi))$$

The Model

$$S_\phi = \int d^4x \sqrt{-g} \left\{ \sum_{j=1}^{j=2} G_j(\phi) X^j + \frac{V}{M_j(\phi) V^{j-1}} \right\}$$

$$G_j(\phi) = \frac{(2\alpha)^{j-1} (1 + f(\phi))^{2-j}}{(1 + f(\phi))^2 + (8\alpha V)}, \qquad M_j = \frac{(1 + f(\phi))^{2j-4} + (8\alpha)^{j-1}}{-}$$

$$X = \frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi$$

Limits of the model

Limit 1

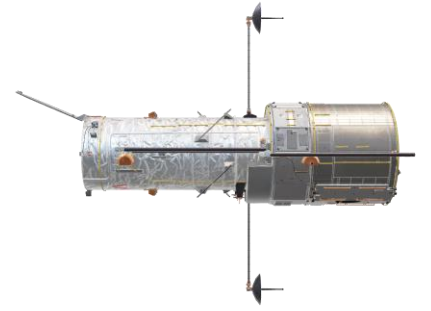
$$2\alpha X \ll (1 + f)$$

Limit 2

$$\bullet \alpha \gg 1$$

Limit 3

$$\bullet \alpha \ll 1$$



slow-roll canonical
inflation

$$S_{limI} = \int d^4x \sqrt{-g} \frac{1}{2} \left\{ R - g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - 2V_{eff}(\phi) \right\}$$

slow-roll K-inflation

$$S_{limII} = \int d^4x \sqrt{-g} \frac{1}{2} \left\{ R + \frac{1}{2} (g^{\mu\nu} \partial_\mu \chi \partial_\nu \chi)^2 - 2V_{eff}(\chi) \right\}$$

Constant-roll
K-inflation

$$S_{limIII} = \int d^4x \sqrt{-g} \frac{1}{2} \left\{ R + \alpha (g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi)^2 - g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - 2V_{eff}(\chi) \right\}$$



Case study

Case study

Case 1

$$V = V_0 \left[1 + \cos \left(\frac{\varphi}{f} \right) \right]$$

$$F(\varphi) = -\xi \varphi^2 R$$

Case 2

$$V = V_0 \left[1 + \cos \left(\frac{\varphi}{f} \right) \right]$$

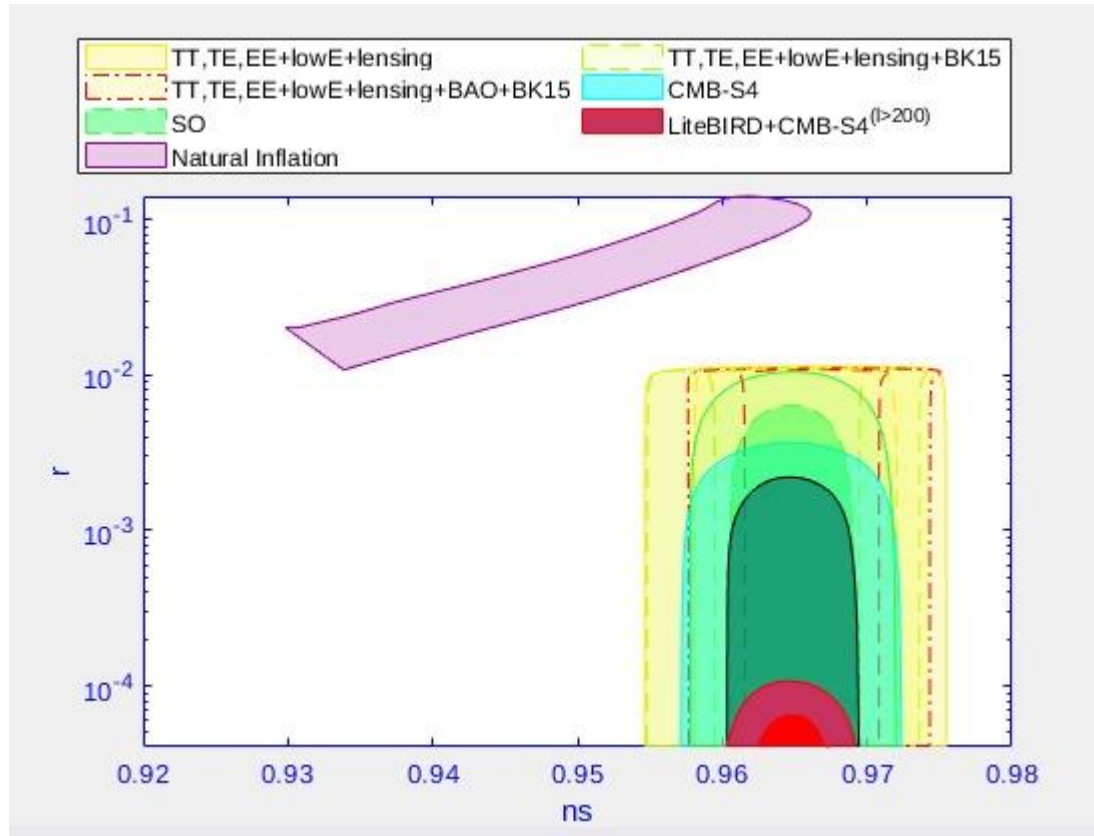
$$F(\varphi) = \lambda \left(1 + \cos \left(\frac{\varphi}{f} \right) \right)$$

Case 3

$$V = V_0 \exp(-l \varphi)$$

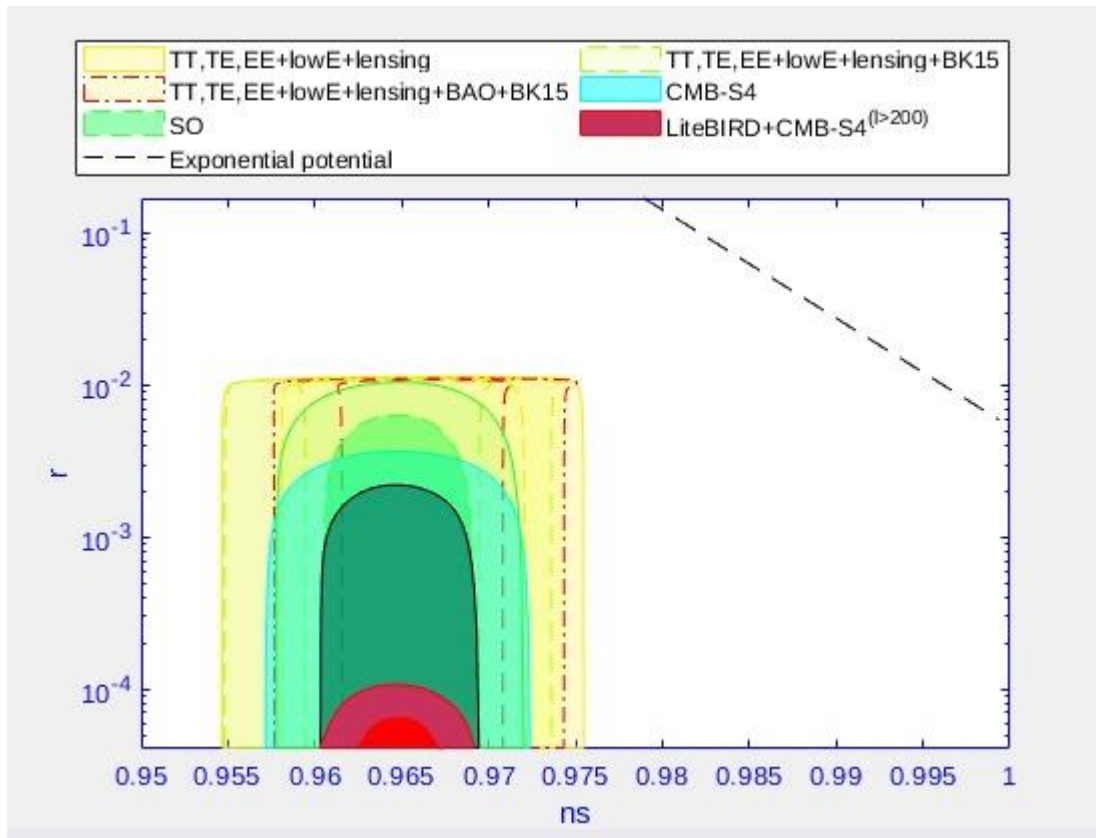
$$F(\varphi) = -\xi \exp(-2l\varphi)$$

Natural Inflation In Standard GR gravity



- The model failed to accommodate the observational results.

Inflation with Exponential potential In Standard GR gravity



- The model failed to accommodate the observational results.

Results

- By studying the effects upon modifying the model, we found that the αR^2 -gravity influences the scalar-tensor ratio r values. In contrast, NMC to gravity significantly impacts the spectral index values (n_s).
- As current observational constraints on parameter α and ξ/λ are pretty loose, having contributions from both terms allows the models to accommodate the observational constraints.

Results: Cosine Potential

M.AIHallak, N.Chamoun and M.S.Eldaher, Natural Inflation with non minimal coupling to gravity in R^2 gravity under the Palatini formalism, JCAP ,2022), 001
 doi:10.1088/1475-7516/2022/10/001

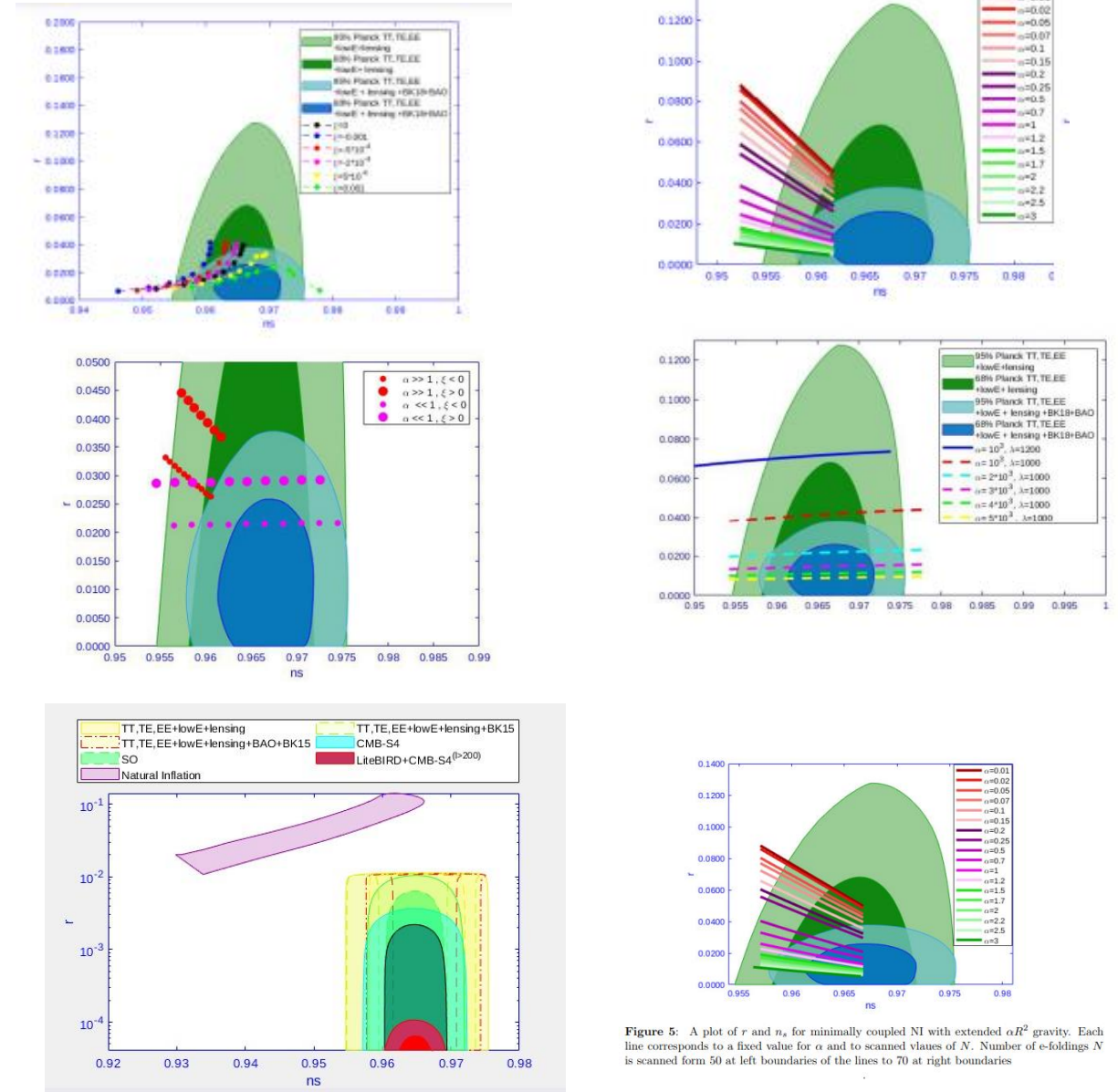
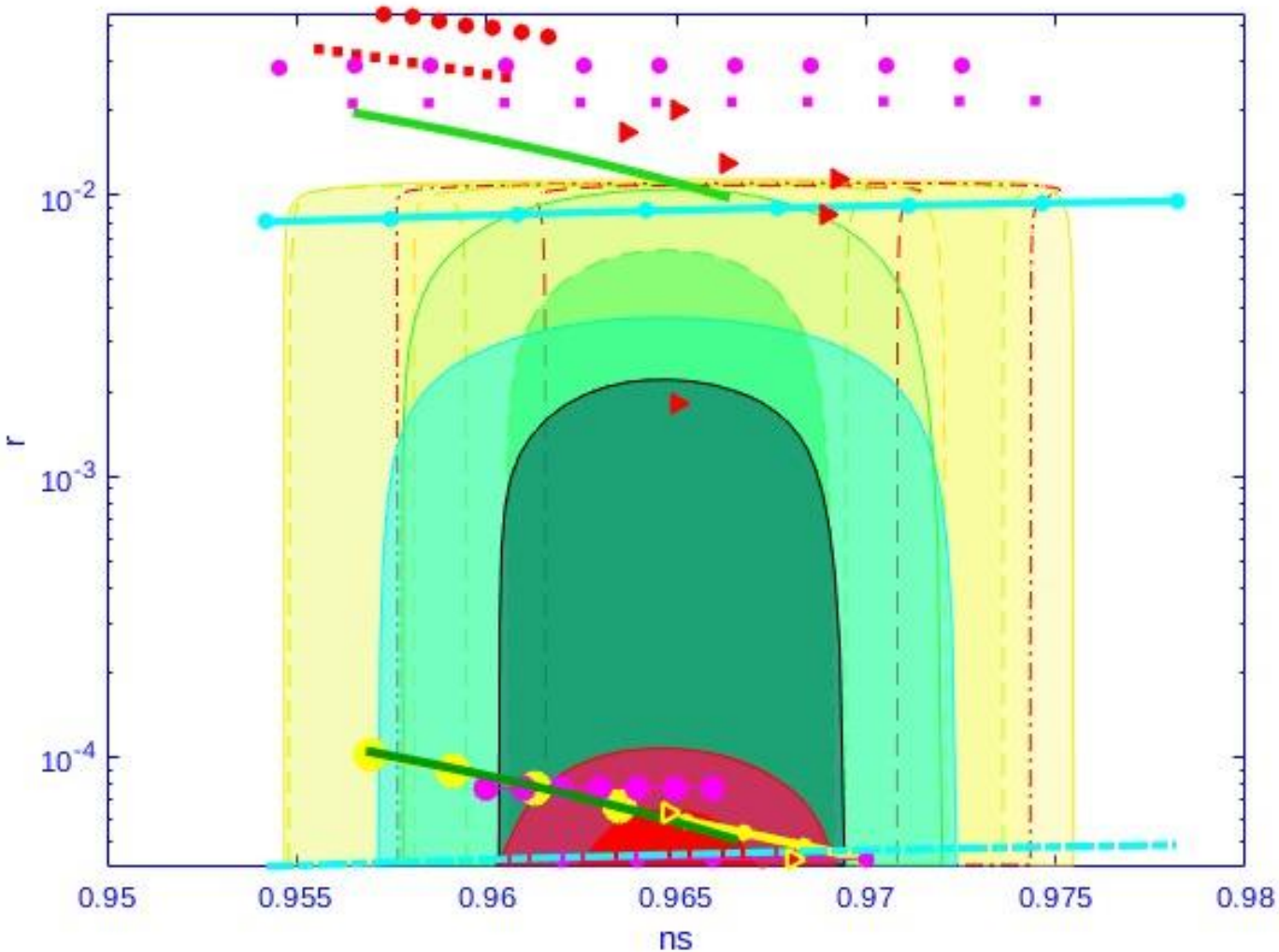
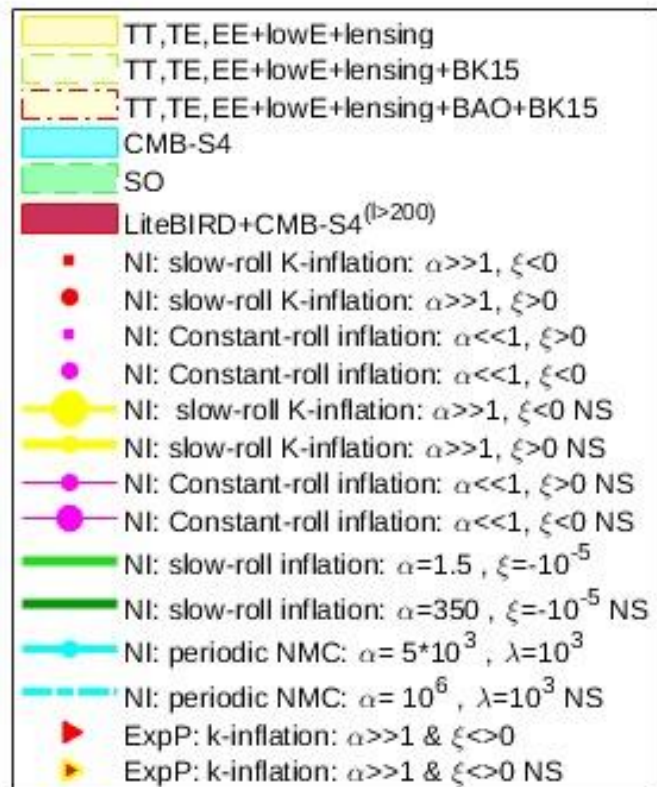


Figure 5: A plot of r and n_s for minimally coupled NI with extended αR^2 gravity. Each line corresponds to a fixed value for α and to scanned values of N . Number of e-foldings N is scanned from 50 at left boundaries of the lines to 70 at right boundaries

Results:



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Thanks for your listening