

# NOETHER SYMMETRIES & INFLATIONARY SCENARIOS IN GHOST-FREE F(R,G) GRAVITY

## 1. INTRODUCTION

Einstein-Gauss-Bonnet models are well-suited for high-curvature regimes relevant to early cosmology, but higher-derivative theories may suffer from **ghost** degrees of freedom due to **Ostrogradsky instabilities**.

Recent **ghost-free extensions of f(R,G) gravity** [1] avoid these issues. However, their complexity often requires reconstruction approaches, where functional forms are imposed externally, neglecting the **intrinsic symmetries**. In this work, viable cosmological scenarios are derived directly from the internal structure of the model using **Noether symmetry methods** and tested for observational viability.

## 2. THEORETICAL FRAMEWORK

$$S = \int d^4x \sqrt{-g} \left[ \frac{R}{2\kappa^2} + \lambda \left( \frac{1}{2} \partial_\mu \chi \partial^\mu \chi + \frac{\mu^4}{2} \right) - \frac{1}{2} \partial_\mu \chi \partial^\mu \chi + h(\chi)G - V(\chi) \right]$$

$\lambda$  - Lagrange multiplier,  $\mu$  - mass-dimension constant,  $\chi$  - auxiliary field

Gauss-Bonnet invariant:  $G = R_{\mu\nu\rho\sigma}R^{\mu\nu\rho\sigma} - 4R_{\mu\nu}R^{\mu\nu} + R^2$        $\chi = \mu^2 t$

$$L = -\frac{3}{\kappa^2} \dot{a}^2 a + a^3 \lambda \left( -\frac{1}{2} \dot{\chi}^2 + \frac{\mu^4}{2} \right) - a^3 \tilde{V} - 8\dot{a}^3 h_\chi \dot{\chi} \quad \tilde{V}(\chi) = -\frac{\mu^4}{2} + V(\chi)$$

## 3. NOETHER SYMMETRY APPROACH

The model contains a conservation law under the condition  $XL = 0$ , where

$$X = \alpha \frac{\partial}{\partial a} + \gamma \frac{\partial}{\partial \chi} + \dot{\alpha} \frac{\partial}{\partial \dot{a}} + \dot{\gamma} \frac{\partial}{\partial \dot{\chi}} \quad \begin{aligned} \dot{\alpha} &= \dot{a} \frac{\partial \alpha}{\partial a} + \dot{\chi} \frac{\partial \alpha}{\partial \chi} \\ \dot{\gamma} &= \dot{a} \frac{\partial \gamma}{\partial a} + \dot{\chi} \frac{\partial \gamma}{\partial \chi} \end{aligned}$$

Conservation law  $Q = \alpha L_{\dot{a}} + \gamma L_{\dot{\chi}}$

The system of equations obtained after dividing the condition  $XL = 0$  in the order of derivatives

$$\begin{aligned} \alpha + 2a\alpha_a &= 0 & \frac{6}{\kappa^2} \alpha_\chi + a^2 \lambda \gamma_a &= 0 & h_\chi \gamma_a &= 0 \\ a\beta + 3\lambda\alpha + 2a\lambda\gamma_\chi &= 0 & h_\chi \alpha_\chi &= 0 \\ 3h_\chi \alpha_a + h_\chi \gamma_\chi + h_{\chi\chi} \gamma &= 0 \\ \frac{1}{2} a \mu^4 \beta + \frac{3}{2} \lambda \mu^4 \alpha - 3V\alpha - aV_\chi \gamma &= 0 \end{aligned}$$

## 4. SOLUTIONS

$$\tilde{V} = C_2 e^{-\frac{3C_1 \chi}{a^{3/2} \gamma_0}}, \quad h = \frac{2\gamma_0 C_3}{3C_1} a^{\frac{3}{2}} e^{\frac{3C_1 \chi}{2a^{3/2} \gamma_0}} + C_4$$

Conservation law

$$Q = -\frac{C_1}{a^{3/2}} \left( \frac{6H}{\kappa^2} + 24H^2 \mu^2 h_\chi \right) - \gamma_0 a^3 (\lambda \mu^2 + 8H^3 h_\chi)$$

These results are used to numerically solve the following equation obtained from a combination of equations of motion

$$\frac{\dot{H}}{\mu^4} + 4H^2 h_{\chi\chi} + \frac{8}{\mu^2} \dot{H} H h_\chi + \frac{3HC_1 a^{\frac{3}{2}}}{\mu^2 \gamma_0} (1 + 4H\mu^2 h_\chi) + \frac{q_0 a^{-3}}{2\mu^2 \gamma_0} = 0$$

## 5. OBSERVATIONAL CONSTRAINTS

