

# Finite-Time Blow-up Solution to a Stochastic Quasilinear Viscolastic Wave Equation with Nonlinear and Logarithmic Source

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## INTRODUCTION & AIM

Stochastic partial differential equations in separable Hilbert spaces have been extensively studied in recent years, with results on existence, uniqueness, stability, blow-up, and other qualitative properties of solutions. In this work, we consider the stochastic viscolastic wave equation with logarithmic nonlinearity:

$$\begin{cases} u_t - \mu \Delta u - (\lambda + \mu) \nabla (dtv) + \int_0^t h(t-s) \Delta u(s) ds \\ + |u|^{p-2} u - u |u|^{p-2} \ln |u|^k + \sigma(x, t) W_t(x, t) & \text{in } \mathcal{D} \times ]0, +\infty[, \\ u(x, t) = 0 & \text{on } \partial \mathcal{D} \times ]0, +\infty[, \\ u(x, 0) = u_0(x), \quad u_t(x, 0) = u_1(x) & \text{in } \mathcal{D}, \end{cases} \quad (1)$$

where:  $\mathcal{D}$  is a bounded domain in  $\mathbb{R}^n$ ,  $n \in \mathbb{N}^*$ , with a smooth boundary  $\partial \mathcal{D}$ ,  $\mu, \lambda$  are the Lamé constants which satisfy  $\mu > 0, \lambda + \mu \geq 0$ ,  $h$  is a positive function, the constant  $k$  is a small non-negative real number, and  $L^2(\mathcal{D})$  is the set of square integrable function on  $\mathcal{D}$  equipped  $W(x, t)$  is an infinite dimensional Wiener process,  $\sigma(x, t)$  is  $L^2(\mathcal{D})$  valued progressively measurable and  $\epsilon$  is positive constant which measures the strength of noise. We rewrite (1) as an equivalent Itô's system

$$\begin{cases} dv - v dt, \\ dv - [\mu \Delta u + (\lambda + \mu) \nabla (dtv) - \int_0^t h(t-s) \Delta u(s) ds \\ - |v|^{p-2} v + u |u|^{p-2} \ln |u|^k] dt + \sigma(x, t) dW_t(x, t) & \text{in } \mathcal{D} \times ]0, +\infty[, \\ u(x, t) = 0 & \text{on } \partial \mathcal{D} \times ]0, +\infty[, \\ u(x, 0) = u_0(x), \quad v(x, 0) = u_1(x) & \text{in } \mathcal{D}, \end{cases} \quad (2)$$

which can be written as the integral equation

$$\begin{cases} u(t) = u_0 + \int_0^t v(s) ds, \\ v(t) = v(0) + \int_0^t [\mu \Delta u + (\lambda + \mu) \nabla (dtv) - \int_0^t h(t-s) \Delta u(s) ds \\ - |v|^{p-2} v + u |u|^{p-2} \ln |u|^k] ds + \int_0^t \sigma(x, s) dW_s(x, t) & \text{in } \mathcal{D} \times ]0, +\infty[, \\ u(x, t) = 0 & \text{on } \partial \mathcal{D} \times ]0, +\infty[, \\ u(x, 0) = u_0(x), \quad v(x, 0) = u_1(x) & \text{in } \mathcal{D}. \end{cases} \quad (3)$$

### Energy Functional

$$\begin{aligned} e(t) &= \frac{1}{2} \|u_t\|_2^2 + \frac{1}{2} \left( \mu - \int_0^t h(s) ds \right) \|\nabla u\|_2^2 + \frac{\lambda + \mu}{2} \|dtv\|_2^2 \\ &+ \frac{1}{2} (h \circ v)(t) + \frac{k}{p^2} \|u\|_p^p - \frac{1}{p} \int_{\mathcal{D}} |u|^p \ln |u|^k dx \end{aligned} \quad (4)$$

$$(h \circ v)(t) = \int_0^t h(t-s) \|v(\cdot, t) - v(\cdot, s)\|_2^2 ds.$$

## METHOD

(A1) Assume that  $h : \mathbb{R}^+ \rightarrow \mathbb{R}^+$  is a  $C^1$  nonincreasing function satisfying

$$h(0) > 0, \quad \mu - \int_0^\infty h(s) ds = l > 0$$

and there exist two nonnegative constants  $\varsigma_1$  and  $\varsigma_2$  such that

$$-\varsigma_1 h(t) \leq h'(t) \leq -\varsigma_2 h(t), \quad t \geq 0.$$

(A2)

$$\int_0^\infty h(s) ds < \mu \frac{(p-2)p}{(p-1)^2}.$$

(A3)  $p > q \geq 2$  and

$$\begin{cases} 2 < p \leq \frac{2(n-1)}{n-2} & \text{if } n \geq 3 \\ 2 < p \leq +\infty & \text{if } n = 1, 2. \end{cases} \quad (5)$$

We purpose

$$G(t) = \frac{\epsilon^2}{2} \sum_{j=1}^{\infty} \mathbb{E} \int_0^t \int_{\mathcal{D}} \lambda_j e_j^2(x) \sigma^2(x, s) dx ds,$$

$$H(t) = G(t) - \mathbb{E}e(t),$$

$$L(t) := H^{1-\alpha}(t) + \delta \mathbb{E}(u, v),$$

where  $0 < \alpha < \min\{\frac{p-1}{2(p+1)}, \frac{p+1-q}{p+1}\}$  and  $\delta$  is a very small constant.

## RESULTS & DISCUSSION

**Theorem 2.1.** Assume (A1) - (A3) hold. Let  $(u, v)$  be a solution of system (2) with initial data  $(u_0, v_0) \in H_0^1(\mathcal{D}) \times L^2(\mathcal{D})$  satisfying

$$\mathbb{E}e(0) \leq -(1 + \beta)E_1, \quad (8)$$

where  $\beta$  is nonnegative constant and  $E_1$  is given (9). If  $p > q$ , then there exists a positive time  $T_0 \in ]0, T[$  such that

$$\lim_{t \rightarrow T_0} \mathbb{E}e(t) = +\infty,$$

where

$$T_0 = \frac{1-\alpha}{\alpha K L^{\frac{\alpha}{1-\alpha}}(0)},$$

$$L(0) = H^{1-\alpha}(0) + \delta \mathbb{E}(u_0, u_1) > 0,$$

and  $K$  is given later.

**proof** We purpose

$$G(\infty) \leq \frac{\epsilon^2}{2} \text{Tr}(Q) \alpha_0^2 \mathbb{E} \int_0^\infty \int_{\mathcal{D}} \sigma^2(x, s) dx ds := E_1 < \infty, \quad (9)$$

where

$$\text{Tr}(Q) = \sum_{j=1}^{\infty} \lambda_j < \infty \quad \text{and} \quad \alpha_0 = \sup_{j \geq 1} \|e_j\|_{\infty} < \infty,$$

where  $Q$  is the covariance operator of Wiener process  $W$ . We showed

$$\begin{aligned} L'(t) &\geq \gamma(H(t) + \mathbb{E}\|v\|_2^2 + \mathbb{E}\|dtv\|_2^2 + \mathbb{E}(h \circ \nabla u)(t) \\ &+ \mathbb{E}\|\nabla u(t)\|_2^2 + \mathbb{E}\|u\|_p^p + \mathbb{E}\|u\|_{p+1}^{p+1}) \\ &\geq 0, \end{aligned} \quad (10)$$

$$\begin{aligned} (L(t))^{1-\alpha} &\leq \tilde{C} \left[ H(t) + \mathbb{E}\|v\|_2^2 + \mathbb{E}\|\nabla u\|_2^2 + \mathbb{E}\|dtv\|_2^2 \right. \\ &\left. + \mathbb{E}\|u\|_p^p + \mathbb{E}(h \circ \nabla u)(t) + \mathbb{E}\|u\|_{p+1}^{p+1} \right], \end{aligned} \quad (11)$$

According to (10) and (11), we have

$$(L(t))^{1-\alpha} \leq \tilde{K} L'(t). \quad (12)$$

In a direct integration of (12), we get

$$(L(t))^{1-\alpha} \geq \frac{1}{(L(0))^{1-\alpha} - \frac{\tilde{K} \alpha t}{1-\alpha}}$$

Therefore,  $L(t)$  blows up in time  $T \leq T_0 = \frac{1-\alpha}{\alpha K L^{\frac{\alpha}{1-\alpha}}(0)}$ .

## CONCLUSION

Under appropriate assumptions on the relaxation function, the damping exponent, and the noise intensity, we construct a suitable energy functional and derive refined energy estimates. By introducing an appropriate Lyapunov functional and exploiting the non-convex structure of the logarithmic potential, we establish sufficient conditions for finite-time blow-up with positive probability. In particular, we prove that if the initial energy is below a critical negative threshold depending explicitly on the noise intensity, then the corresponding solution cannot exist globally in time. Moreover, an explicit upper bound for the blow-up time is obtained in terms of the initial data and the system parameters.

## FUTURE WORK / REFERENCES

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