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2 **A proposal for introducing quantum physics in the** 3 **footsteps of Einstein**

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10 **Abstract:** We formulate a didactic proposal for introducing some fundamental concepts of quantum
11 physics to advanced high school students, and to their teachers. The inspiration comes from some
12 of the fundamental papers about the subject by Albert Einstein, in which many of these concepts,
13 for example light quanta, wave-particle duality and probability, were introduced for the first time,
14 in a characteristically illuminating way. The proposal can be supplemented by a discussion of
15 elementary tools of statistical physics, which are needed at some point. Preliminary results, both
16 with students and teachers, are very promising.

17 **Keywords:** Physics teaching; history of physics; quantum mechanics

19 **1. Introduction**

20 Nowadays physics teachers and researchers are greatly challenged by the need to develop
21 educational tracks aimed at introducing advanced high school students to the main concepts of
22 quantum theory [1], in order to engage them with modern physics and to stimulate their interest
23 towards science and technology. This is made more urgent by the fact that many high school teachers
24 may lack a proper advanced physics education (for example if their major was in mathematics, as it
25 often happens in Italy). Hence, teachers themselves should benefit from appropriately designed
26 educational paths. It has been repeatedly emphasized that the history of physics can play an
27 important role in teaching, in multiple ways (see e.g. [2]). One of the possibilities, especially for
28 modern physics, is taking inspiration from the reading of original texts by the founders of the subject.
29 This is especially the case for quantum physics, the teaching of which is very often developed by
30 following a historical path. Indeed, the typical teaching-learning sequence for quantum physics starts
31 from Planck's formula for black body radiation, and gradually introduces – roughly in the same order
32 in which they were discovered – the main physical concepts that in the end will be subsumed in the
33 full theory of quantum mechanics. Usual textbook treatments of elementary quantum physics employ
34 a rather standardized choice of topics, of course at different depths, depending on their target, be it
35 undergraduate or high school students. Examples include [3,4,5,6]. While we have no objection
36 towards this way of developing the subject, we recognize that history is far more complex and rich
37 than any didactic presentation. While it is of course natural that textbooks which are not history
38 manuals simplify the historical development, it may still be the case that some additional insight
39 could be gained by looking at usually neglected bits of history. In the present paper we suggest that
40 such insight can be found in some of Einstein's fundamental works in quantum physics. Actually, as
41 is well known by historians, many of the fundamental concepts that nowadays are at the core of
42 quantum mechanics, have been introduced in some paper by Einstein [7]. Despite this, most of these

43 works are often ignored in teaching. For example, after discussing Planck’s radiation formula and its
44 justification in terms of energy quantization, it is customary to go on immediately to the explanation
45 of the photoelectric effect in terms of light quanta. This contribution, revolutionary as it is, is often
46 the only one by Einstein that is mentioned. The deep difference between Planck’s hypothesis of
47 energy quantization and Einstein’s light quanta is often not emphasized. Nor is there any mention of
48 Einstein’s statistical reasoning in order to arrive at the latter [8]. And the great conceptual leap
49 involved in going from the energy quantization or light quanta hypotheses, which are statistical in
50 nature, to the application to the photoelectric effect, in which *individual quanta* are involved, is often
51 not stressed enough. After this, traditional courses go straight to discussing Bohr’s atomic model, and
52 then to De Broglie’s hypothesis of matter waves. The pivotal introduction by Einstein of the concept
53 of wave-particle duality for light (later extended to matter), again by a statistical reasoning applied
54 to the black body radiation formula [9], which shows that both the wave and the particle nature of
55 light are necessary for accounting for the observed spectrum, is usually neglected. Going on, while
56 Einstein’s modern derivation of the Planck formula in terms of atomic transitions [10] is sometimes
57 treated (especially in view of its important application to lasers), the really revolutionary part of it,
58 namely the intrinsic probabilistic and causality violating nature of the process of spontaneous
59 emission is not properly emphasized. Yet, this was the first appearance of probably the most mind-
60 blowing characteristics of quantum mechanics. Notable exceptions in the literature, which do discuss
61 one of more of these aspects, include [11,12], as well as explicitly historically oriented works like [13,
62 14]. But these are quite advanced works, while typical high school books completely neglect these
63 developments. However, we believe that also advanced high school students should benefit from a
64 discussion of the above mentioned topics, if of course they are presented in a suitably simple way.
65 The same is true for teachers, with the advantage of their wider mathematical and physical
66 background, which allows a deeper treatment. An important feature of many of Einstein’s papers is
67 that they are uniformly models of clarity and clear physical thinking. Also, unlike many papers by
68 Einstein’s contemporaries, these works turn out to be very readable and enlightening also for modern
69 readers. Moreover, all of Einstein’s papers up to 1927 (for now) are freely available in English
70 translation of the web [15], and translations in many other languages are often republished. These
71 papers, full of marvellous insight, deep philosophical discussions and examples of how great science
72 is done, can therefore constitute a wonderful and stimulating reading for high school teachers and
73 also, in a suitable selection, for pupils.

74 Inspired by the above considerations, we have extracted a didactic path to quantum physics
75 from the papers by Einstein we have cited. We have tried to stay as close as possible to Einstein’s
76 original reasonings, while trying to simplify the mathematics whenever possible, so that nothing
77 more complicated than elementary integral calculus is used, but without losing physical insight. This
78 makes the derivations affordable by advanced high school students. Moreover, modern notation is
79 used throughout. Emphasis is on clarity rather than on mathematical rigor.

80 **2. The track**

81 In this section we give some details on the didactic track we developed. In order to be concise,
82 those parts of it which rely on standard material which is readily available in the literature are only
83 sketched. It is assumed that students have studied the standard curriculum of classical physics, up to
84 electromagnetic waves, and including the kinetic theory of gases. Moreover, the ability to perform
85 simple integrals is taken for granted. Students in the second half of their last year of high school,
86 should therefore be prepared to undertake this path.

87 *2.1. Step 1: Cavity radiation and Planck’s law*

88 As most teaching-learning sequences in quantum physics, our one begins with Planck’s law for
89 the energy density of radiation in a cavity at thermal equilibrium:

90

$$u(\nu, T) = \frac{8\pi h\nu^3}{c^3} \frac{1}{e^{\frac{h\nu}{k_B T}} - 1}. \quad (1)$$

91

92 This law can be introduced in a phenomenological way as a function which fits experimental data on
93 thermal radiation (this is actually how Planck first introduced it), and it can be treated in a standard
94 way. A heuristic explanation of why energy quantization allows to get Planck's formula can be given
95 in terms of how the former affects the principle of equipartition of energy, which when applied to
96 electromagnetic waves would give the unacceptable Rayleigh-Jeans (RJ for short) law. Several aspects
97 need to be emphasized, in order for what follows to be clearer. In particular, the universality of black
98 body radiation, and hence of the constants which appear in it, specifically Planck's constant h . Then
99 it is important to see how the limit $\frac{h\nu}{k_B T} \ll 1$, in which energy quanta are much smaller than the
100 average thermal energy, which is called the classical limit (in this limit h cancels out), gives back the
101 Rayleigh-Jeans law, while the opposite limit $\frac{h\nu}{k_B T} \gg 1$, which may be dubbed the "extremely quantum
102 limit", gives a limiting form of the Planck law, which is the Wien law:
103

$$u(\nu, T) = \frac{8\pi h\nu^3}{c^3} e^{-\frac{h\nu}{k_B T}}. \quad (2)$$

104 This is an important step since Einstein introduced light quanta by studying the latter law,
105 which is valid precisely in the regime where the quantum nature of light is most evident.

106 2.2. Step 2: Einstein 1905: light quanta

107 In his 1905 paper [8], Einstein computes in an ingenious way the entropy density associated with
108 thermal radiation described by Wien's distribution (2). The computation is quite easy, since it
109 involves standard thermodynamics and the integration of a logarithmic function. The result can be
110 used to obtain the entropy variation under an isothermal and adiabatic expansion from a volume V_0
111 to a volume V . This is given by:

$$S(V) - S(V_0) = k_B \ln \left(\frac{V}{V_0} \right)^{\frac{E}{h\nu}}, \quad (3)$$

112 and is valid separately for each frequency component of the radiation. This can be compared with
113 the entropy variation associated with the Joule expansion (that is, a free expansion which is both
114 adiabatic and isothermal) of the ideal gas, which is as well known $S(V) - S(V_0) = k_B \ln \left(\frac{V}{V_0} \right)^N$, where
115 N is the number of molecules making up the gas. This expression can be easily and instructively
116 derived from Boltzmann's statistical formula for entropy. It is striking that the two expressions are
117 exactly the same, provided one identifies
118

$$\frac{E}{h\nu} = N \quad \text{or} \quad E = N h\nu. \quad (4)$$

119 Einstein's bold interpretation of the above identification is that radiation behaves statistically as if it
120 were made up of N independent "quanta", each of which has energy given by Planck's expression $h\nu$,
121 so that the total energy of the radiation is just the sum of the energy of individual quanta, which
122 behave just like the molecules of an ideal gas. It is important to emphasize that this statement is much
123 more revolutionary than Planck's hypothesis, according to which only the exchanges of energy
124 between the radiation and the walls of the cavity are discrete. Here instead it is argued that the
125 radiation *itself* is quantized, at least in the limit where quantum effects dominate and Wien's law is

126 valid. In doing this, Einstein for the first time used statistical physics in a peculiar and illuminating
127 way. In fact, he used the expression of the entropy coming from thermodynamics to infer the
128 statistical behaviour of radiation, while usually the opposite is done.

129 2.3. Step 3: Some applications

130 At this point it is useful to consider some applications of the above result. These applications
131 actually go one step further because assume that not only the statistical behavior of the radiation is
132 given by its corpuscular nature, but that the physical effects of *individual* quanta can be measured.
133 Hence they show that quantization of radiation is a fundamental feature of Nature. In his paper [8],
134 Einstein considers three applications, one of which is the famous one to the photoelectric effect, which
135 of course must be included in the track. It can be treated in a standard way (e.g. as in [3]). We have
136 found that also another of these applications, which is usually not treated, namely the explanation of
137 Stokes' rule for fluorescence, can be useful and instructive. Such rule, in fact, states simply that, when
138 a fluorescent material absorbs light, it reemits it with a lower frequency. A well-known effect is the
139 violet glow of white clothes which are exposed to ultraviolet light¹. This rule is readily
140 understandable in terms of light quanta. Namely, fluorescent materials absorb light quanta, whose
141 energy is $\varepsilon = h\nu$, and reemit them. Since part of the energy is absorbed, light quanta are reemitted
142 with a lower energy $\varepsilon' < \varepsilon$. But $\varepsilon' = h\nu'$, hence $\nu' < \nu$, which is precisely Stokes' rule. The third
143 application concerns ionization by UV light quanta, and it can be considered as well.

144 2.4. Step 4. Einstein 1909: wave-particle duality

145 Going back to the full Planck distribution, an interesting consideration can be made. In the
146 classical limit, it reduces to the RJ distribution, which can be obtained by applying equipartition of
147 energy to electromagnetic waves. On the other hand, in the "most quantum" limit, it reduces to the
148 Wien distribution, which as we just saw leads to considering radiation as an ideal gas of non-
149 interacting quanta. But the full Planck distribution contains both these limits. This means that, while
150 at extreme values of the frequencies an entirely wave or an entirely particle picture works well, for
151 the full distribution both aspects are necessary. Giving a quantitative foundation to this observation
152 was what Einstein did in 1909 [9]. In that paper, Einstein brilliantly applied a statistical reasoning to
153 the full Planck radiation law (1). This time, he focused on quadratic fluctuations of the energy, which
154 can be computed by using elementary statistical physics methods, namely Boltzmann's distribution
155 of energy. At this point a slight detour may be advisable, since this topic is not usually known by
156 students. The result of the computation, which once again does not involve difficult mathematical
157 steps, is

158

$$(\Delta E)^2 = (h\nu u + \frac{c^3 u^2}{8\pi \nu^2})Vd\nu. \quad (5)$$

159 Again, this result is separately valid for all frequencies. The first term has exactly the same form of
160 what one could get from an ideal gas, when one restricts to molecules having all the same energy
161 $\varepsilon = h\nu$ (the analog of monochromatic radiation), which is $(\Delta E)^2 = \varepsilon^2 N$, where $N = E/h\nu$ is the
162 average number of molecules., and it is the only term one gets when one considers the Wien
163 distribution in place of the Planck one. The other term is what one would obtain by considering waves
164 (the fluctuations are quadratic in the energy density due to interference effects), and it is the only
165 contribution if the RJ law is used. Eq. (5) shows that the Planck distribution involves both a particle
166 contribution and a wave contribution. Einstein's genius revealed itself in the choice of the right
167 quantity, in which the two contribution are unentangled, and appear as a simple sum. This result,

¹ This is usually well known by students, especially those who like going to the disco.

168 Einstein concluded, was leading to (in his words) “a theory of light which can be interpreted as a kind of
169 fusion of the wave and the emission [i.e. the corpuscular] theory.

170 2.5. Step 5. Interlude: the Bohr atom

171 Einstein’s next breakthrough paper, which we are discussing in the next subsection, was inspired
172 by Bohr’s atomic model. Therefore this model has to be properly treated. Again, this can be done in
173 the usual way [3]. The important result for what follows is that, according to that model, electrons in
174 atoms and molecules (not only in the hydrogen atom) can only occupy a set of stationary states Z_1 ,
175 Z_2 , etc. with discrete energies E_1 , E_2 , etc., and transitions between these states can occur. According
176 to the Bohr theory postulates, in the transition between the m -th state and the n -th state, with $E_m >$
177 E_n , if the electron goes from Z_m to Z_n , it emits radiation with frequency $E_m - E_n = h\nu$, while the
178 inverse transition can be induced if the electron absorbs radiation with the same frequency. A point
179 which is not often emphasized is that, when he proposed his model, Bohr did not believe in light
180 quanta, hence he did not phrase his postulates in terms of these entities.

181 2.6. Step 6. Einstein 1917: probability

182 In 1916 [10], Einstein went back to the problem of cavity radiation, this time using insight coming
183 from the Bohr theory. Unlike Bohr, Einstein explicitly assumed that in any transition, a single light
184 quantum having energy equal to the energy difference is emitted or absorbed. Einstein considered
185 radiation in thermal equilibrium with an ensemble of atoms or molecules, which therefore
186 continuously absorb or emit light quanta. In fact, three kinds of elementary processes can occur,
187 namely absorption, stimulated emission (in which a light quantum hits an atom, stimulating its decay
188 to a lower energy state with emission of another quantum), and spontaneous emission (in which an
189 atom spontaneously goes to a state with lower energy emitting a quantum). Since the system is
190 assumed to be at thermal equilibrium, the various states of the atoms will be distributed according
191 to the Boltzmann distribution, and moreover this distribution should not change with time, meaning
192 that on average each transition from a state to another must happen with the same rate as the inverse
193 transition. This translates in a detailed balance equation, from which the Planck distribution for the
194 energy density of radiation follows. Two very important remarks are in order. First, the process of
195 spontaneous emission is very different from the other two, since it is not triggered by the interaction
196 with a light quantum. Instead, it seems to have no apparent cause, and it is necessary to assume that
197 it occurs, in a given interval of time, with a given *probability*. Nor the direction in which the emerging
198 light quantum is emitted is predictable. That is, spontaneous emission is an intrinsically probabilistic
199 process, which seems to violate the principle of causality. This behavior is actually analogous to that
200 of radioactive decay of a nucleus. Therefore, the Planck distribution can only be reproduced if
201 intrinsically probabilistic and causality violating processes occur. This means that the quantum
202 behaviour is inextricably linked with probability, a fact which is of course included in the full
203 quantum mechanics, but which had not been anticipated.

204 3. Results

205 The above track was tested both with teachers and (in a reduced form for the moment) with
206 selected students, and in the course of various outreach activities performed at the University of
207 Salerno. The preliminary results are very encouraging. Teachers’ responses also emphasized how they
208 managed to see quantum physics from a new point of view, which helped them to overcome some
209 criticalities in their understanding. Remarkably, also those whose major subject at university was physics
210 shared this impression, since they did not encounter this way of treating the subject in their studies. The
211 students found the material intriguing, not least because of the association with Einstein’s name, and
212 they developed some intuition about the quantum phenomena which were the object of study.

213 4. Discussion

214 Since basically all of Einstein's fundamental ideas on light quanta came from statistical
215 mechanics, of which he was a true master, it is possible to enhance our path by introducing the
216 necessary tools of this subject, starting from the elementary notions of kinetic theory that are part of
217 the standard curriculum. In particular, a knowledge of Boltzmann's entropy formula for the first part,
218 and of Boltzmann's distribution of energies for the second part, is necessary. These tools can be
219 introduced by again taking inspiration from Einstein's treatment, when available (in fact, he often
220 added some review of the statistical tools he needed in his papers). Teaching elementary statistical
221 physics in high schools is however a very active area of research on its own, with a large body of
222 literature (see e.g. [16, 17]), hence our proposal can also act as a link between the latter and the
223 teaching of quantum physics. The possibility of using some of Einstein's papers in this endeavor as
224 well has been the subject of a parallel study to the present one, and will be reported elsewhere.

225 Of course, the path developed in this paper does not constitute a complete course in quantum
226 physics, for several reasons. First of all, relevant work by others still needs to be considered of course.
227 Moreover, the path entirely unfolds in the realm of the old quantum theory. After it is completed, an
228 introduction to the basics of the full theory of quantum mechanics must be given. Also for this, a very
229 rich literature has developed (see e.g. [1] and references therein), hence there are various ways in
230 which our proposal can be complemented. We have developed our own proposal in [18], again taking
231 some inspiration from the heroic history of quantum mechanics. Preliminary tests of an enlarged
232 didactic path involving a combination of both the proposals are beginning soon.

233 5. Conclusions

234 In this paper, we described a possible didactic path for elementary quantum physics, which, as
235 part of the literature on the teaching of quantum physics was inspired by the history of physics. The
236 preliminary results we obtained indicate that this path can usefully complement already existing
237 material, both in forming high school teachers, and in directly teaching students.

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239 investigate the pedagogical value of Einstein's 1905 paper.

240 **Author Contributions:** The paper is the result of discussions between the three authors. M.D. began to test parts
241 of this proposals on students and teachers in the course of outreach activities performed at the university of
242 Salerno. He and A.N. analyzed the responses of students and teachers. M.D. wrote the paper.

243 **Conflicts of Interest:** The authors declare no conflict of interest.

244 Abbreviations

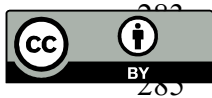
245 The following abbreviations are used in this manuscript:

246 RJ: Rayleigh-Jeans

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